# Microscopic collective inertial masses in nuclear reaction

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Low-energy nuclear reaction – Macroscopic model

$$\left\{-\frac{1}{2\mu_R}\frac{d^2}{dR^2} + \frac{L(L+1)}{2\mu_R R^2} + V(R)\right\}\psi_L(R) = E_L\psi_L(R)$$

-How good/bad is this?

2020.5.18「第3回クラスター階層領域研究会」 via Zoom

### Nuclear reaction with shape evolution

- Low-energy nuclear reaction
  - Relative motion between two nuclei
  - Shape change



Is the relative coordinate *R* meaningful?

### Nuclear reaction with shape evolution

- One-to-one correspondence
  - *R* can be a choice to define the scale of the coordinate.
  - R affects the inertial masses



Relative distance:  $M(R) \sim \mu_R \quad (R \to \infty)$ Orientation:  $I(R) \sim 2\mu_R R^2 \quad (R \to \infty)$ 

### Inertial mass

 $\mathcal{X}$ 

X

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• A particle moving along the *x* axis

$$-L = \frac{1}{2}m\dot{x}^2$$

• Assuming the motion along the X axis

$$-L = \frac{1}{2}m\dot{X}^{2}$$
 (Wrong)

• Representing in the X axis (x = f(X))

$$-L = \frac{1}{2}m_{eff}\dot{X}^2$$
 (Correct)

$$-m_{eff} = \frac{m}{(\cos\theta)^2}$$

# Effect of "effective mass"

- Velocity-dependent potential
- Nucleonic effective mass

$$-\frac{m^*}{m} \sim 0.7 - 0.8$$

• Does this affect the inertial mass of nuclear reaction?

 $-(M(R), I(R)) \to (\mu_R, 2\mu_R R^2) \times \frac{m^*}{m}?, \text{ at } R \to \infty$ 

### Construction of macroscopic model

Model Hamiltonian

$$\left\{-\frac{d}{dR}\frac{1}{2M(R)}\frac{d}{dR} + \frac{L(L+1)}{2I(R)} + V(R)\right\}\psi_L(R) = E_L\psi_L(R)$$

Microscopically calculating V(R), M(R), I(R)

Can we reproduce the following asymptotic values at  $R \rightarrow \infty$ ? How good is the usage of these values?

$$M(R) = \mu_R, \qquad I(R) = \mu_R R^2$$

# ASCC method

- Optimal reaction path based on TDHF dynamics
- Inertial masses with residual effect beyond mean fields (cf. Thouless-Valatin)
  - Neglecting the residual effect

– <u>Cranking formula</u> for collective masses

$$M_{\rm cr}^{\rm NP}(R) = 2 \sum_{n \in p, j \in h} \frac{|\langle \varphi_n(R) | \partial / \partial R | \varphi_j(R) \rangle|^2}{e_n(R) - e_j(R)},$$

#### 3D real space representation

- 3D space discretized in lattice
- BKN functional:  $E_{\text{BKN}}[\rho, \tau]$  (rather schematic)
- Moving mean-field eq.: Imaginary-time method
- Moving RPA eq.: Finite amplitude method (PRC 76, 024318 (2007))



At a moment, no pairing

1-dimensional reaction path extracted from the Hilbert space of dimension of  $10^4 \sim 10^5$ .



Different masses between ASCC & cranking after touching

 $\square$ 

Reaction path deviates from R



"Cranking mass" decreases with decreasing  $m^*$ 

$$M(R) \neq \mu_R \quad (R \to \infty)$$



### Moment of inertia: I(R)



I(R) decreases with reducing  $m^*$ 

# Moment of inertia



Smooth transition from  $I_{rigid}$  to  $I_{red}$ .

### Quantum effect

Quantum spherical systems: I = 0 (no rotation)

$$m_{1}$$

$$I(R) = I_{red}(R) + I_{1} + I_{2}$$

$$I_{red}(R) = \mu_{R}R^{2} = \frac{m_{1}m_{2}}{m_{1} + m_{2}}R^{2}$$

$$I_{i}: Mol of nucleus i w.r.t. its CoM$$

$$m_{1} + m_{2}$$

$$I(R) = I_{red}(R) + I_{rigid}^{(1)} + I_{rigid}^{(2)}$$

# Alpha reaction: $^{16}O + \alpha$

Synthesis of <sup>20</sup>Ne

#### **Fusion reaction: Astrophysical S-factor**

$$\sigma(E) = \frac{1}{E} P(E) \times S(E)$$

Effect of dynamical change of the inertial mass

1 1.5 2 2.  
$$E_{c.m.}$$
 [MeV]





# Summary

- Self-consistent description of nuclear reaction path and dynamics
  - Reaction path affects the inertial mass
  - Effective mass is canceled by the residual effect in the asymptotic region (the cranking formula fails)
  - Vanishing MoI in spherical nuclei (quantum effect) is properly taken into account
  - Reduction of astrophysical S-factor

K. Wen, T. N., Front. Phys. 8 (2020) 16; Phys. Rev. C 96 (2017) 014610