

3-body force & beyond in 1D cold atoms

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第8回クラスター階層領域研究会

Feb.9 - Feb.11 (2023) @ Osaka

1. 3-body force naturally arises in 1D

- **Thermal conductivity**

T. Tanaka & Y. Nishida, *Phys. Rev. E* 106, 064104 (2022)

“Thermal conductivity of a weakly interacting Bose gas in quasi-one-dimension”

- **Cluster formation**

Y. Nishida, *Phys. Rev. A* 97, 061603(R) (2018)

“Universal bound states of 1-dimensional bosons with 2- and 3-body attractions”

2. Artificial control of 3-body force & beyond

- **Quantum droplet**

Y. Sekino & Y. Nishida, *Phys. Rev. A* 97, 011602(R) (2018)

“Quantum droplet of one-dimensional bosons with a three-body attraction”

- **Efimov effect**

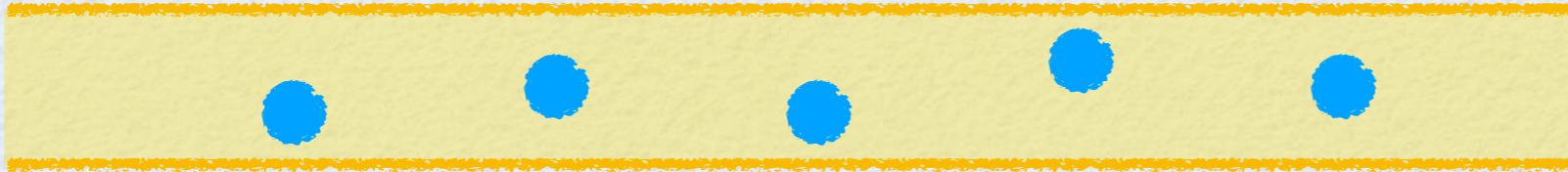
Y. Nishida & D. T. Son, *Phys. Rev. A* 82, 043606 (2010)

“Universal four-component Fermi gas in one dimension”

3-body force in 1D

Bosons in quasi-1D

Confining 3D bosons into a tight 1D waveguide

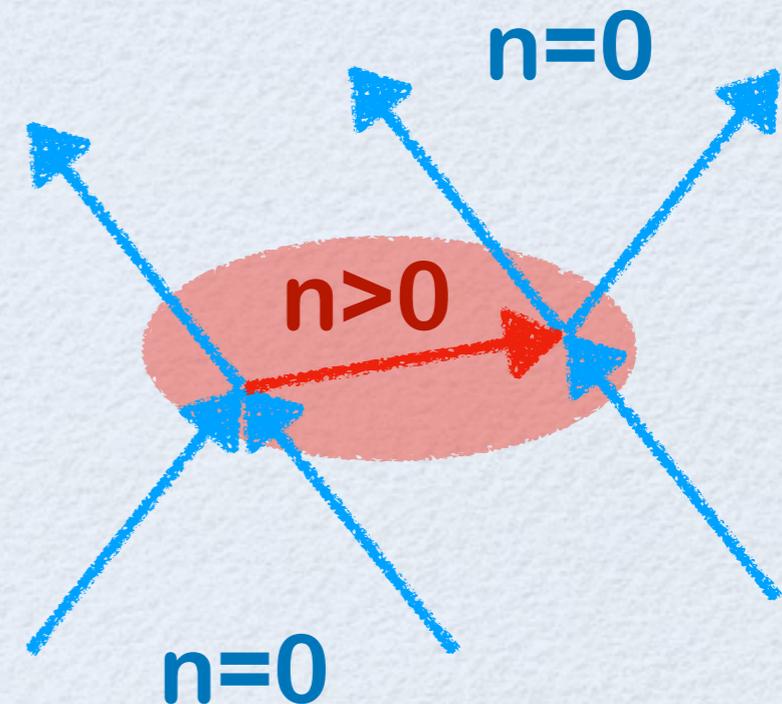


$$\hat{\mathcal{H}}_{3\text{D}} = \hat{\Phi}^\dagger \left(-\frac{\nabla^2}{2m} + \frac{x^2 + y^2}{2ml_\perp^4} \right) \hat{\Phi} + \frac{2\pi a_{3\text{D}}}{m} (\hat{\Phi}^\dagger \hat{\Phi})^2$$



$$T, \mu \ll \hbar\omega_\perp, \quad |a_{3\text{D}}| \ll l_\perp$$

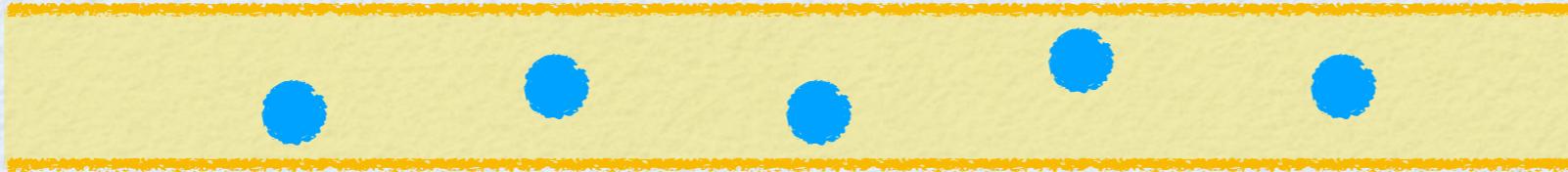
$$\hat{\mathcal{H}}_{1\text{D}} = -\hat{\phi}^\dagger \frac{\partial_x^2}{2m} \hat{\phi} + \frac{g_2}{2} (\hat{\phi}^\dagger \hat{\phi})^2$$



effective 3-body interaction

I. E. Mazets & J. Schmiedmayer, NJP (2010)

Confining 3D bosons into a tight 1D waveguide



$$\hat{\mathcal{H}}_{3\text{D}} = \hat{\Phi}^\dagger \left(-\frac{\nabla^2}{2m} + \frac{x^2 + y^2}{2ml_\perp^4} \right) \hat{\Phi} + \frac{2\pi a_{3\text{D}}}{m} (\hat{\Phi}^\dagger \hat{\Phi})^2$$



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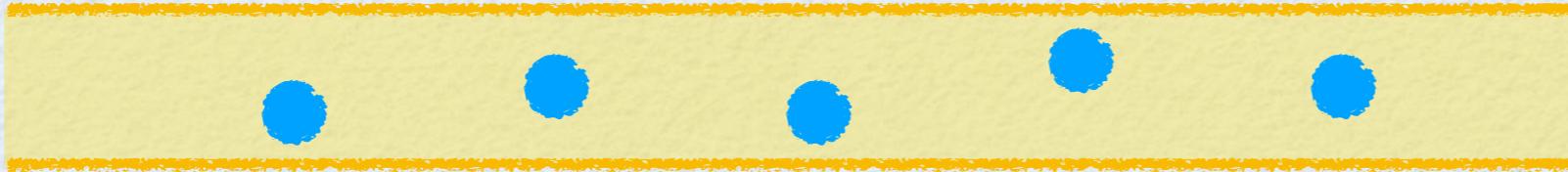
$$\hat{\mathcal{H}}_{1\text{D}} = -\hat{\phi}^\dagger \frac{\partial_x^2}{2m} \hat{\phi} + \frac{g_2}{2} (\hat{\phi}^\dagger \hat{\phi})^2 + \frac{g_3}{6} (\hat{\phi}^\dagger \hat{\phi})^3 + \dots$$

$$\text{with } g_2 = 2 \frac{a_{3\text{D}}}{ml_\perp^2}, \quad g_3 = -12 \ln(4/3) \frac{a_{3\text{D}}^2}{ml_\perp^2}$$

effective 3-body interaction

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$$\hat{\mathcal{H}}_{3\text{D}} = \hat{\Phi}^\dagger \left(-\frac{\nabla^2}{2m} + \frac{x^2 + y^2}{2ml_\perp^4} \right) \hat{\Phi} + \frac{2\pi a_{3\text{D}}}{m} (\hat{\Phi}^\dagger \hat{\Phi})^2$$



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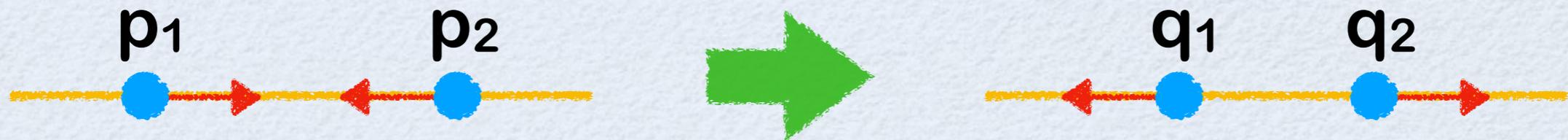
$$\hat{\mathcal{H}}_{1\text{D}} = -\hat{\phi}^\dagger \frac{\partial_x^2}{2m} \hat{\phi} + \frac{g_2}{2} (\hat{\phi}^\dagger \hat{\phi})^2 + \frac{g_3}{6} (\hat{\phi}^\dagger \hat{\phi})^3 + \dots$$

integrability
preserving

integrability
breaking

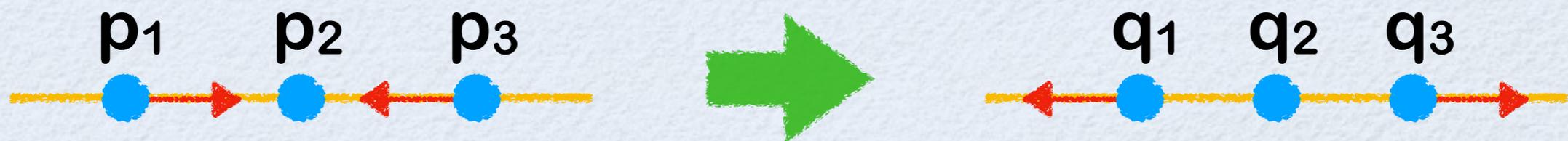
$$\hat{\mathcal{H}}_{1D} = -\hat{\phi}^\dagger \frac{\partial^2_x}{2m} \hat{\phi} + \frac{g_2}{2} (\hat{\phi}^\dagger \hat{\phi})^2 + \frac{g_3}{6} (\hat{\phi}^\dagger \hat{\phi})^3 + \dots$$

2-body scattering $\Rightarrow \{p_1, p_2\} = \{q_1, q_2\}$



cannot change momentum distribution
due to energy and momentum conservations

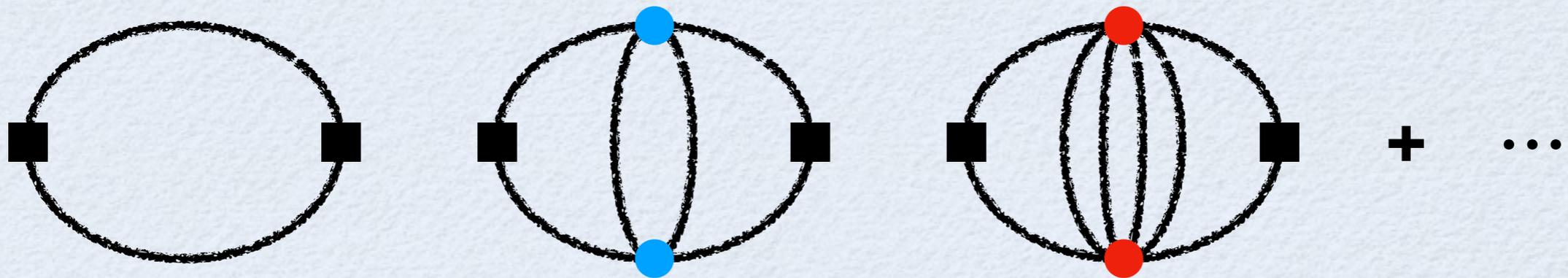
3-body scattering $\Rightarrow \{p_1, p_2, p_3\} \neq \{q_1, q_2, q_3\}$



can change momentum distribution

$$\hat{\mathcal{H}}_{1D} = -\hat{\phi}^\dagger \frac{\partial_x^2}{2m} \hat{\phi} + \frac{g_2}{2} (\hat{\phi}^\dagger \hat{\phi})^2 + \frac{g_3}{6} (\hat{\phi}^\dagger \hat{\phi})^3 + \dots$$

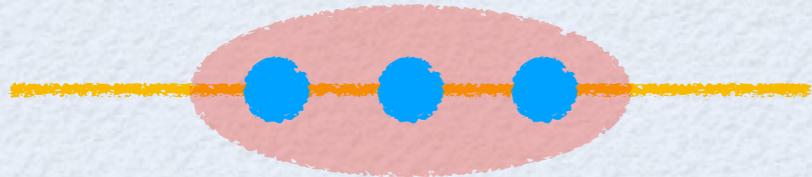
Thermal conductivity at weak coupling is evaluated by Kubo formula + resummation



➔ $\kappa = \frac{\mathcal{N}}{m^3 g_3^2} \tilde{\kappa}(\tilde{T}) = \frac{(l_\perp / a_{3D})^4 \mathcal{N}}{[12 \ln(4/3)]^2 m} \tilde{\kappa}(\tilde{T})$

2-body coupling cancels out
and 3-body coupling dominates

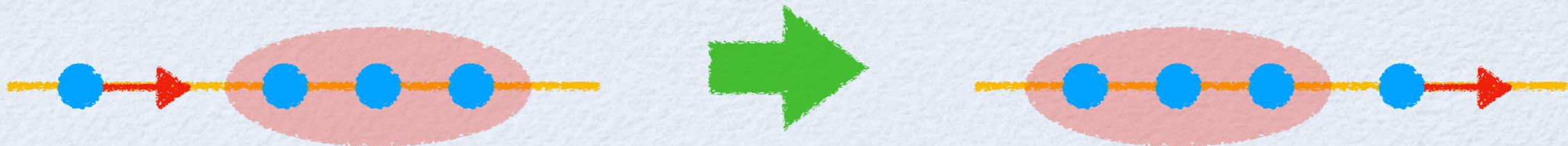
N bosons with **attractive** g_2 form a cluster



$$E_N = -\frac{N(N^2 - 1)}{24}mg_2^2$$

J. B. McGuire, J. Math. Phys. (1964)

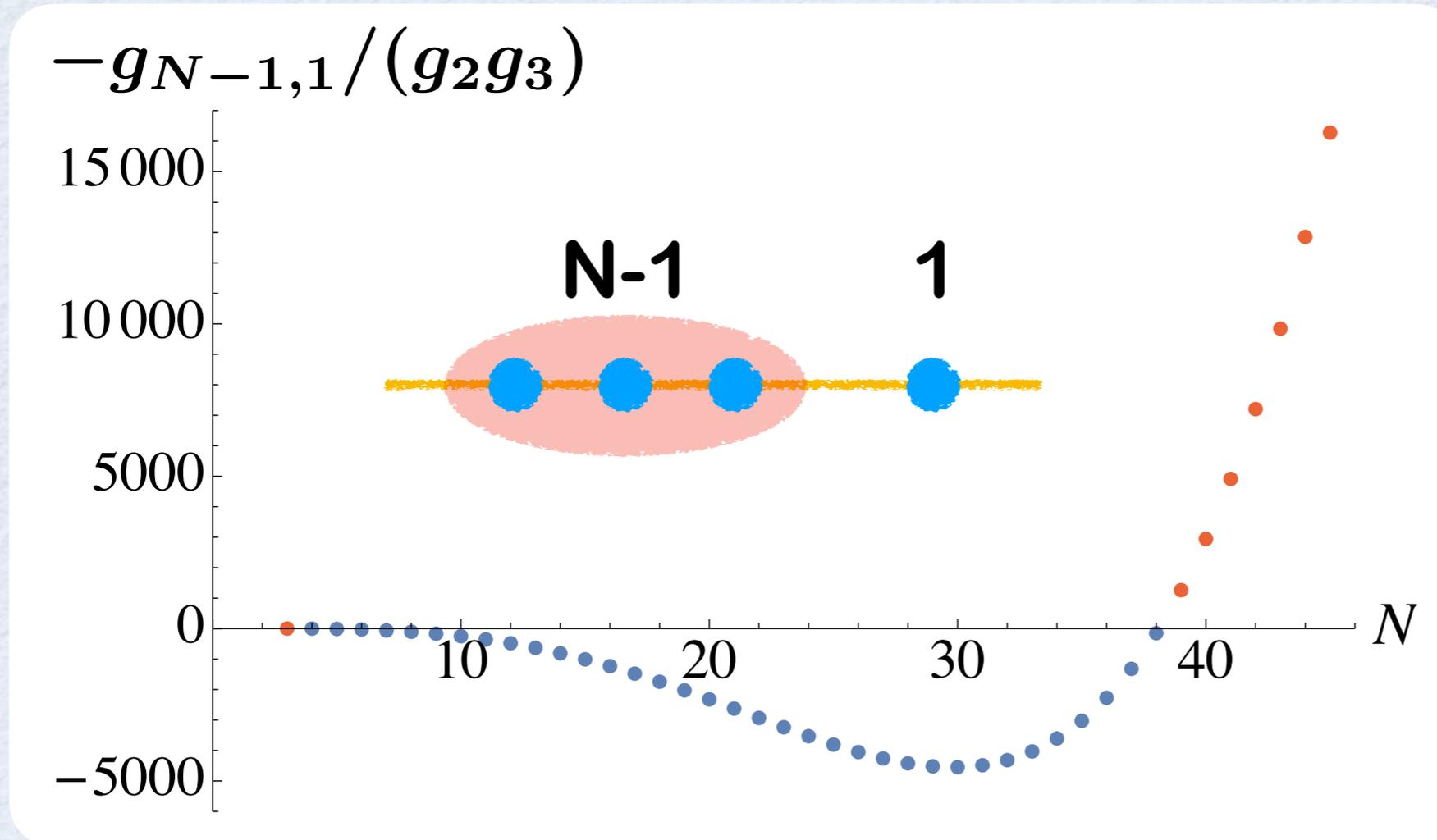
N-body cluster + extra boson



No interaction

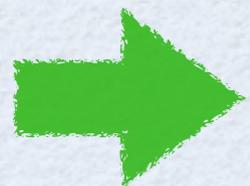
Interaction between cluster & boson is
also dominated by **3-body coupling**

Cluster-boson coupling $g_{N-1,1}$ induced by g_3

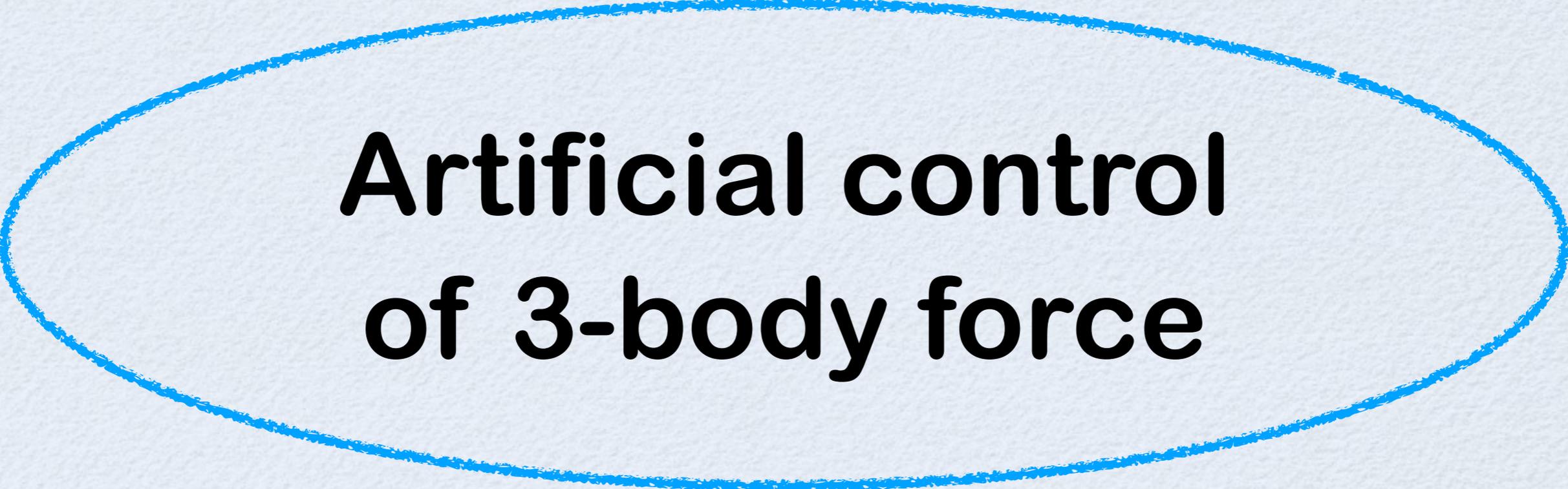


Y. N., PRA (2018)

- Coupling is **repulsive** for $N=4,5,\dots,38$
- Coupling is **attractive** for $N=3$ & $39,40,\dots$



New N-body cluster formation by **3-body force**



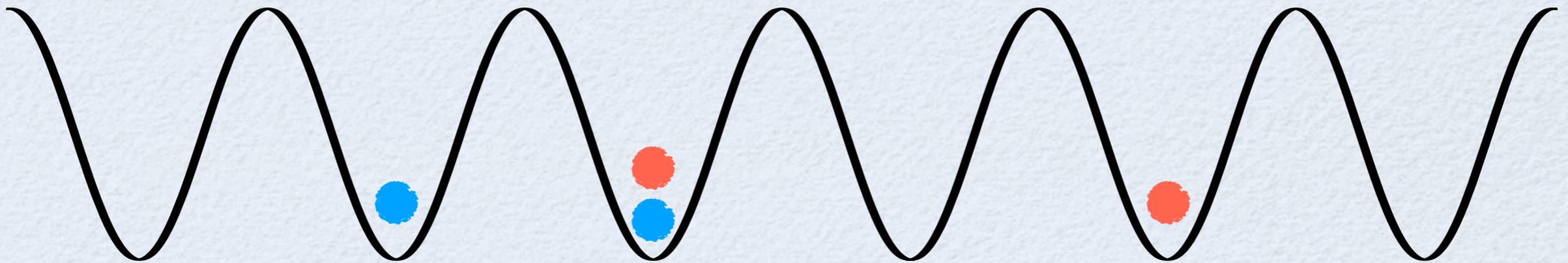
**Artificial control
of 3-body force**

Tunable 2-body & 3-body force

12/18

Coupled 2-component bosons in an optical lattice

$$\hat{H} = -t \sum_{\langle i,j \rangle} \sum_{\sigma} \hat{b}_{\sigma i}^{\dagger} \hat{b}_{\sigma j} + \sum_i \sum_{\sigma, \tau} \frac{g_{\sigma\tau}}{2} \hat{b}_{\sigma i}^{\dagger} \hat{b}_{\tau i}^{\dagger} \hat{b}_{\tau i} \hat{b}_{\sigma i} \quad \text{D. S. Petrov} \\ \text{PRA (2014)}$$
$$- \frac{\Delta}{2} \sum_i (\hat{b}_{\uparrow i}^{\dagger} \hat{b}_{\uparrow i} - \hat{b}_{\downarrow i}^{\dagger} \hat{b}_{\downarrow i}) - \frac{\Omega}{2} \sum_i (\hat{b}_{\uparrow i}^{\dagger} \hat{b}_{\downarrow i} + \hat{b}_{\downarrow i}^{\dagger} \hat{b}_{\uparrow i})$$



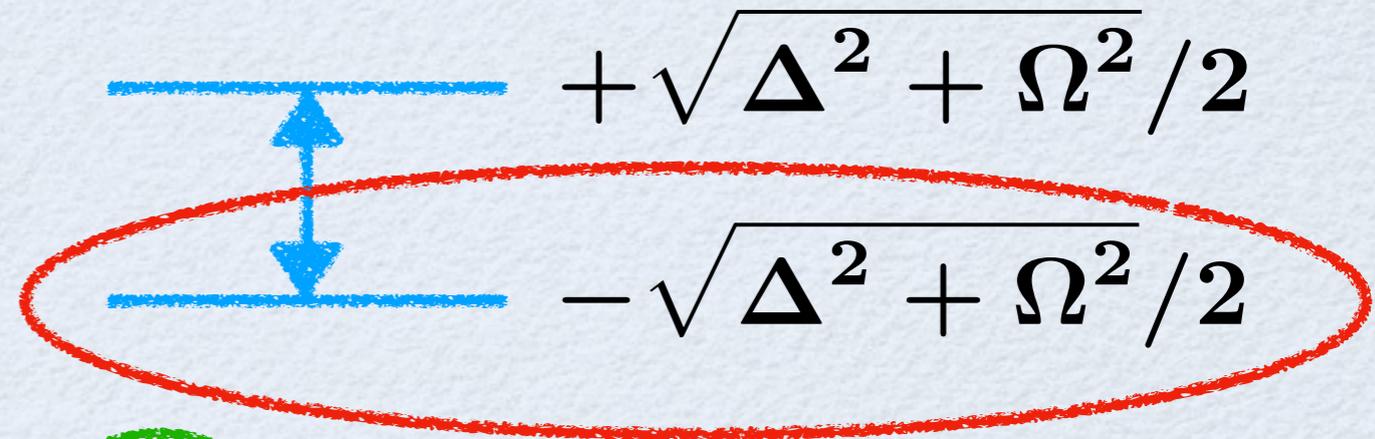
Tunable 2-body & 3-body force

Coupled 2-component bosons in an optical lattice

$$\hat{H} = -t \sum_{\langle i,j \rangle} \sum_{\sigma} \hat{b}_{\sigma i}^{\dagger} \hat{b}_{\sigma j} + \sum_i \sum_{\sigma, \tau} \frac{g_{\sigma\tau}}{2} \hat{b}_{\sigma i}^{\dagger} \hat{b}_{\tau i}^{\dagger} \hat{b}_{\tau i} \hat{b}_{\sigma i} \quad \text{D. S. Petrov PRA (2014)}$$
$$- \frac{\Delta}{2} \sum_i (\hat{b}_{\uparrow i}^{\dagger} \hat{b}_{\uparrow i} - \hat{b}_{\downarrow i}^{\dagger} \hat{b}_{\downarrow i}) - \frac{\Omega}{2} \sum_i (\hat{b}_{\uparrow i}^{\dagger} \hat{b}_{\downarrow i} + \hat{b}_{\downarrow i}^{\dagger} \hat{b}_{\uparrow i})$$



$$T, \mu \ll \sqrt{\Delta^2 + \Omega^2}$$

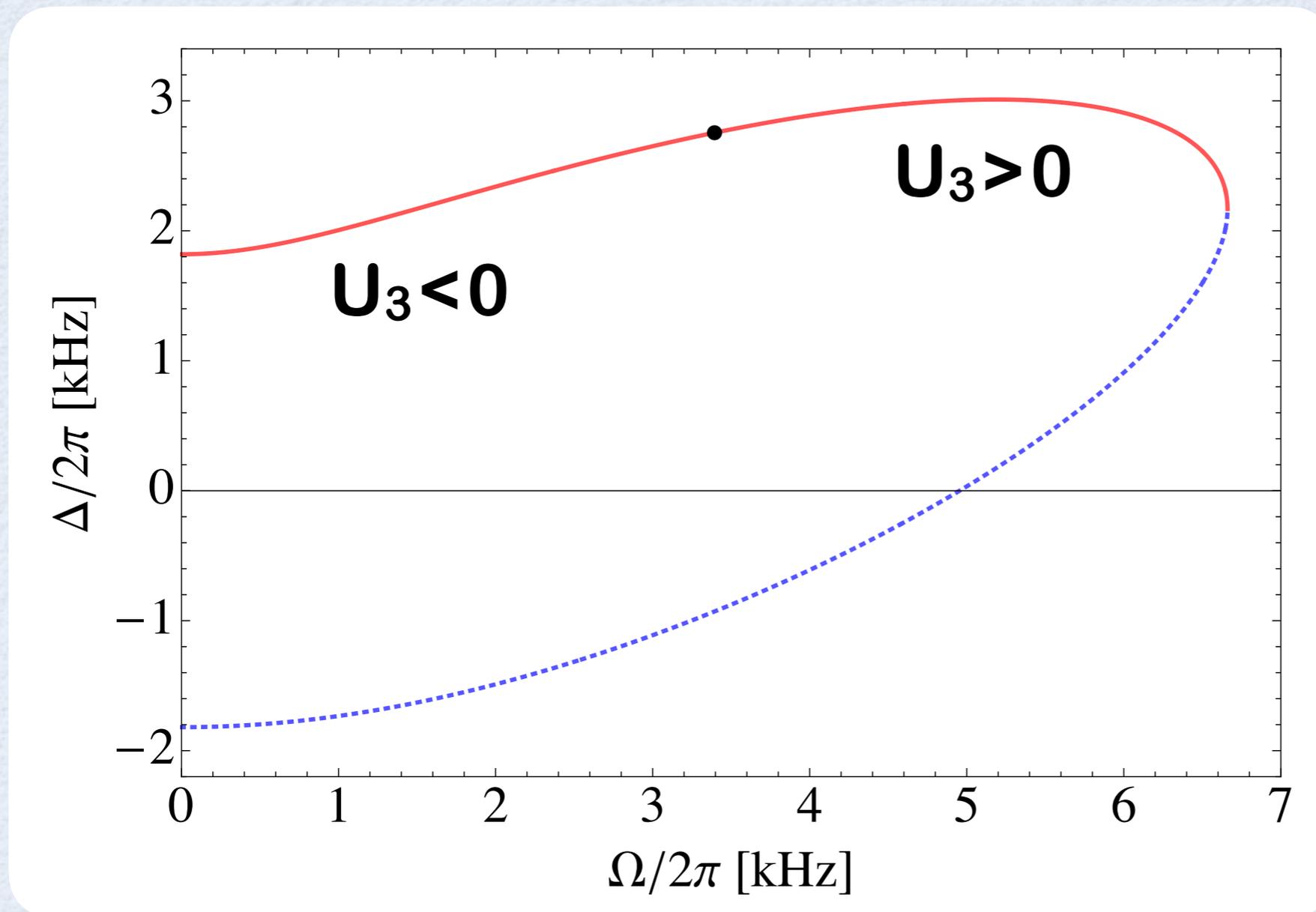


$$\hat{H}_{\text{eff}} = -t \sum_{\langle i,j \rangle} \hat{b}_i^{\dagger} \hat{b}_j + \sum_i \sum_{N=1}^{\infty} \frac{U_N}{N!} (\hat{b}_i^{\dagger} \hat{b}_i)^N$$

Effective N-body coupling depends on Δ , Ω , $g_{\sigma\tau}$

Tunable 2-body & 3-body force

Independent control of 2-body & 3-body coupling
is possible by tuning Δ , Ω



Y. Sekino & Y. N.
PRA (2018)

Line of $U_2=0$ for ^{39}K @ $B=58\text{G}$

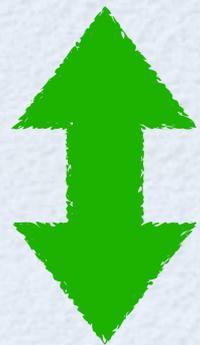
Tunable 2-body & 3-body force

1D bosons with N-body interaction ONLY

$$\hat{\mathcal{H}} = -\hat{\phi}^\dagger \frac{\partial_x^2}{2m} \hat{\phi} + \frac{g_N}{N!} (\hat{\phi}^\dagger \hat{\phi})^N$$

N-body problem with N-body interaction in 1D

⇒ N-1 relative coordinates



$$\hat{H}_N = -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial r_1^2} + \cdots + \frac{\partial^2}{\partial r_{N-1}^2} \right) + g_N \delta^{N-1}(\vec{r})$$

2-body problem with 2-body interaction in D dimensions

⇒ D relative coordinates

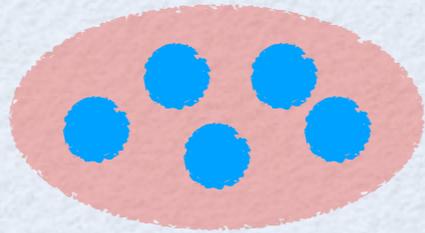
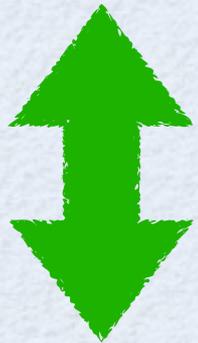


N-body interaction in 1D “realizes”

2-body interaction in (N-1) dimensions

$N \gg 1$ bosons with 2-body attraction in 2D

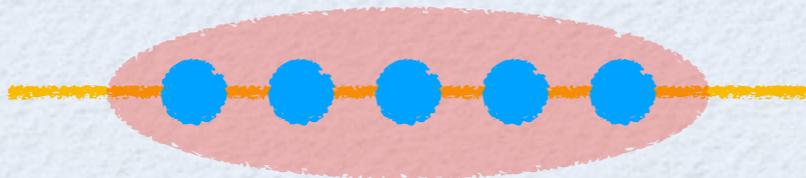
form a **quantum droplet** with $E_N \sim \exp(2.15N)$



H.-W. Hammer & D. T. Son, PRL (2004)

$N \gg 1$ bosons with 3-body attraction in 1D

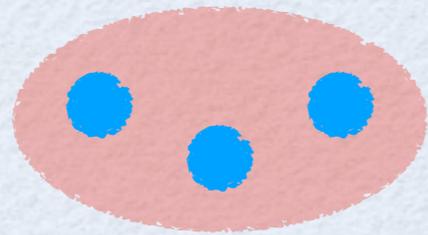
form a **quantum droplet** with $E_N \sim \exp(1.47N^2)$



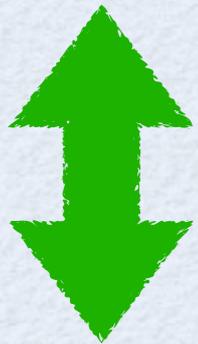
Y. Sekino & Y. N., PRA (2018)

3 bosons at 2-body resonance in 3D

show **Efimov effect** with $E_n \sim (22.7)^{-2n}$



V. Efimov, PLB (1970)

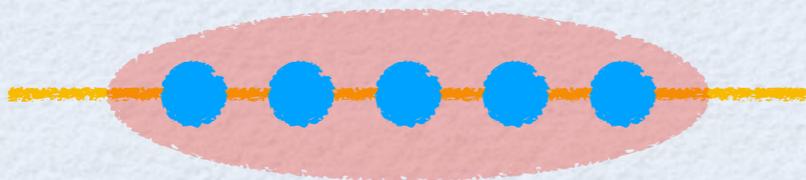


Efimov effect is usually possible only in 3D

E. Nielsen, et al., Phys. Rep. (2001)

5 bosons at 4-body resonance in 1D

show **Efimov effect** with $E_n \sim (12.4)^{-2n}$



D. T. Son & Y. N., PRA (2010)

Unusual realization of Efimov effect in 1D !

1. 3-body force naturally arises in 1D

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T. Tanaka & Y. Nishida, *Phys. Rev. E* 106, 064104 (2022)

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Y. Sekino & Y. Nishida, *Phys. Rev. A* 97, 011602(R) (2018)

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